

# Pauliprinzip & viele Elektronen

Schrödingergl. für  $\Psi(\vec{r}_1, \vec{s}_1; \vec{r}_2, \vec{s}_2; \dots; \vec{r}_N, \vec{s}_N)$  lösen.

(Terme für  $E_{kin}$ ,  $E_{kern-e}$ ,  $E_{ee}$ )

hoffnungslos

effektive Einteilchen S.Gl mit eff. e-e-Potential  $U_{ee}$   
(mean field  $\rightarrow$  s.FH)

$$U_{ee}(\vec{r}) = -e \int d\vec{r}' \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} \quad \text{für } e \text{ bei } \vec{r}$$

$$\text{mit } \rho(\vec{r}) = -e \sum_i |\psi_i(\vec{r})|^2$$

Hartree-Gleichung

Es fehlt das Pauliprinzip!

Deshalb statt einfachem Produktansatz

$$\Psi_1(\vec{r}_1, \vec{s}_1) \cdot \Psi_2(\vec{r}_2, \vec{s}_2) \cdot \dots \cdot \Psi_N(\vec{r}_N, \vec{s}_N)$$

eine Slaterdeterminante als Ansatz:

$$\Psi(\vec{r}_1, \vec{s}_1, \dots) = \begin{vmatrix} \Psi_1(\vec{r}_1, \vec{s}_1) & \Psi_1(\vec{r}_2, \vec{s}_2) & \dots & \Psi_1(\vec{r}_N, \vec{s}_N) \\ \Psi_2(\vec{r}_1, \vec{s}_1) & & & \vdots \\ \vdots & & & \vdots \\ \Psi_N(\vec{r}_1, \vec{s}_1) & & & \Psi_N(\vec{r}_N, \vec{s}_N) \end{vmatrix}$$

⇒ Hartree-Fock-Gleichung

ok. für Atome & kleine Moleküle

für  $F_k$  i.A. nicht lösbar.

außer für FEG.

$$\text{Resultat: } \frac{E}{N} = \frac{e^2}{2a_0} \left[ \underbrace{\frac{3}{5} (k_F a_0)^2}_{\text{mittl. } E_{\text{kin}}} - \underbrace{\frac{3}{2\pi} k_F a_0}_{E_x} \right]$$

( $\nabla$  HF überschätzt Austausch, weil Abschirmung fehlt)

• bei geringen Dichten: Wignerkristall)

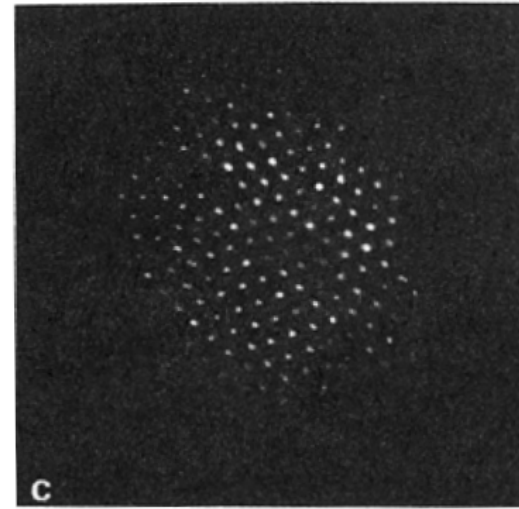
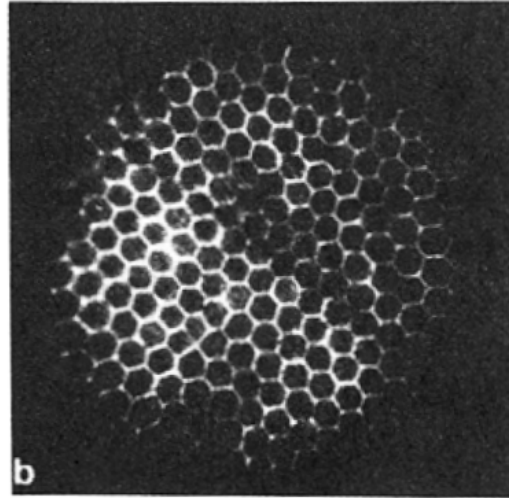
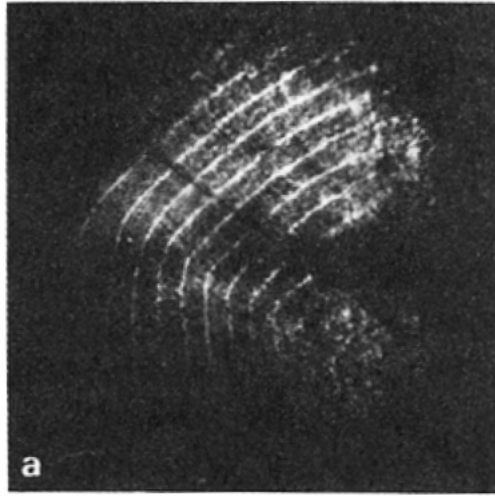
E. Wigner, On the Interaction of Electrons in Metals, Phys. Rev. **46**, 1002 (1934)

|, If the electrons had no kinetic energy, they would settle in configurations which correspond to the absolute minima of the potential energy.

These are closed-packed lattice configurations, with energies very near to that of the bodycentered lattice. |,

*Wie verringert man  $E_{kin}$ ?*

# Elektronen auf flüssigem He



aus: *Electrons at the Surface of Quantum Systems*,  
P. Leiderer, *J. Low Temp. Phys.* 87(1992)

Fig. 11. Formation of a dimple lattice on a  ${}^4\text{He}$  surface ( $T=3.5\text{ K}$ ) charged with electrons from above. The pictures show the surface deformation approximately  $2s$ (a) and  $6s$ (b) and (c) after the field had been increased to  $E_c$ . The image plane in (a) and (b) was chosen such that convex deformations of the surface, corresponding to local maxima, appear bright; in (c) bright areas correspond to local minima (i.e. the center of the dimples). The distance between adjacent rows of dimples is close to the wavelength  $2\pi a$  of the soft ripplon,  $0.24\text{ nm}$  in this case.

aktueller: 2DEGs in (Al)GaAs-Schichtsystemen:

$n$  mit Potentialen einstellbar