

Fig. 4.8. One-dimensional semi-infinite lattice model potential

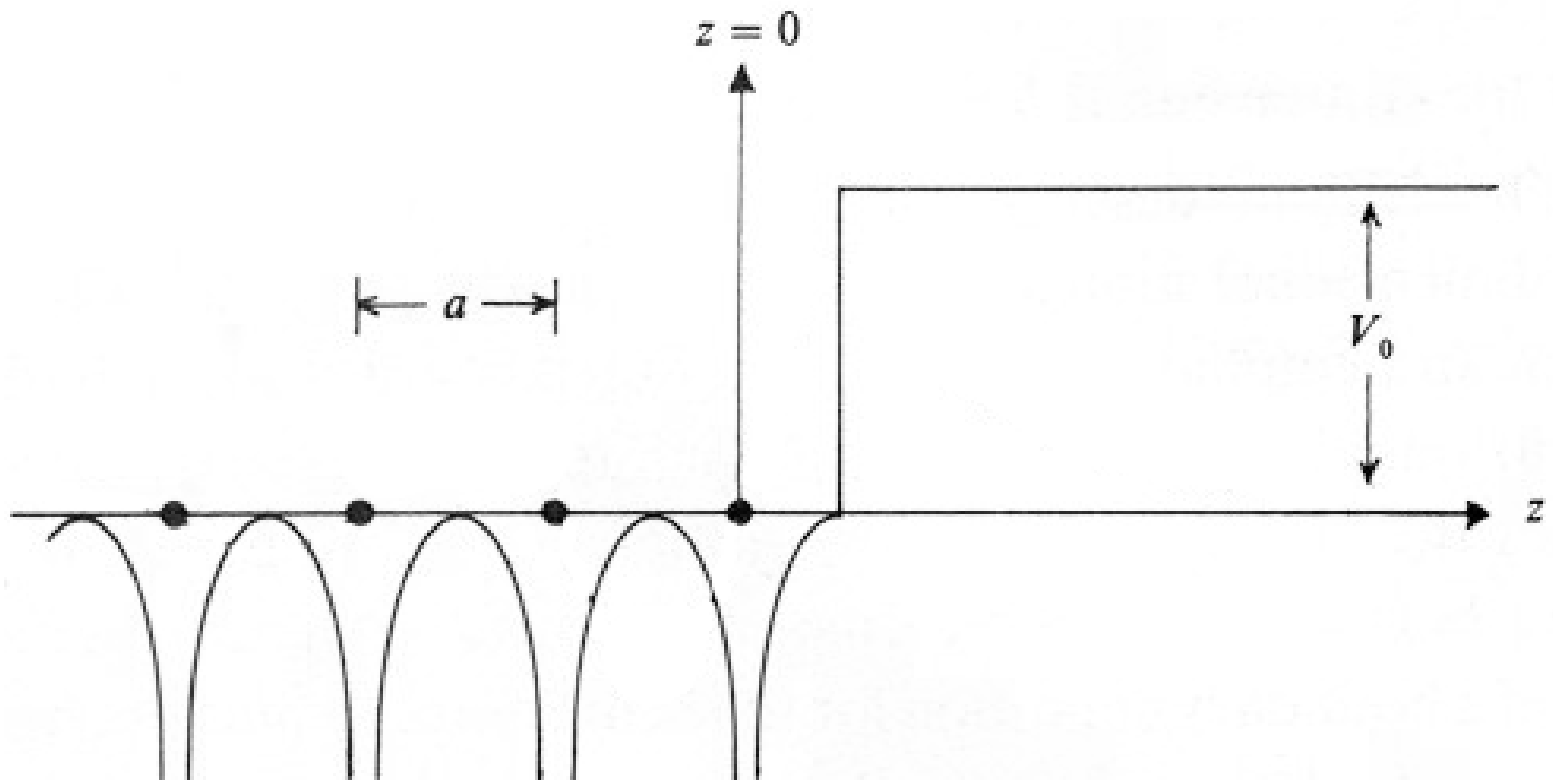
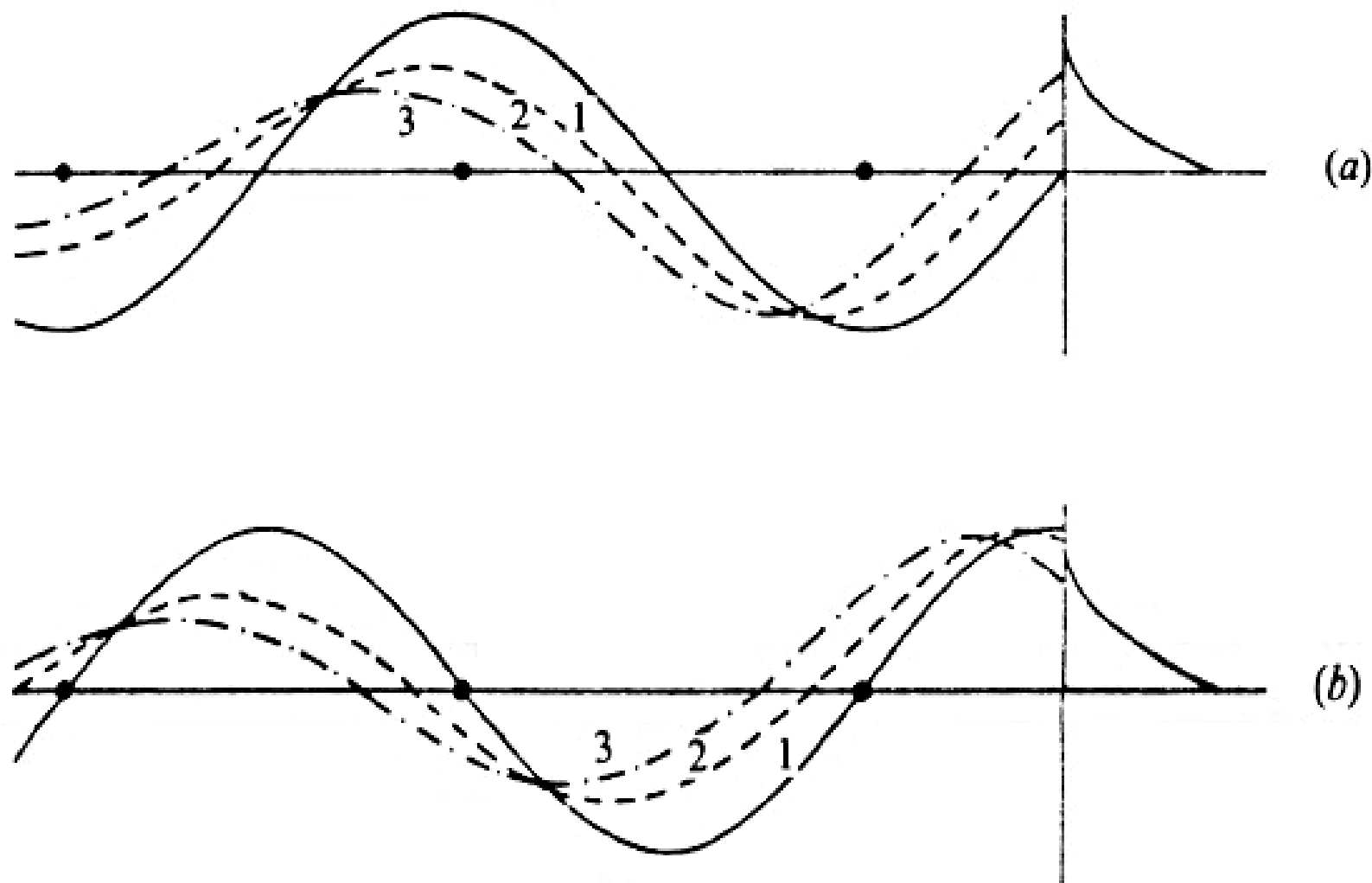


Fig. 4.10. Wave function matching at  $z = a/2$ : (a)  $V_g < 0$ ; (b)  $V_g > 0$ . The sequence 1, 2, 3, indicates increasing energy starting from the bottom of the gap (Forstmann, 1970).



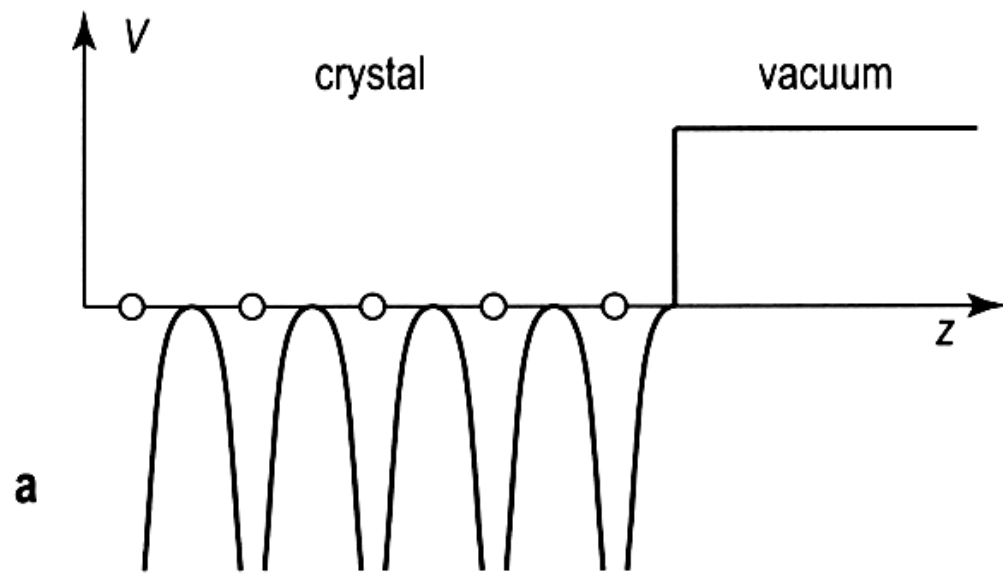


Fig. 11.5. (a) One-dimensional semi-infinite lattice model potential. Two types of wave functions of a semi-infinite crystal: (b) bulk states and (c) surface states

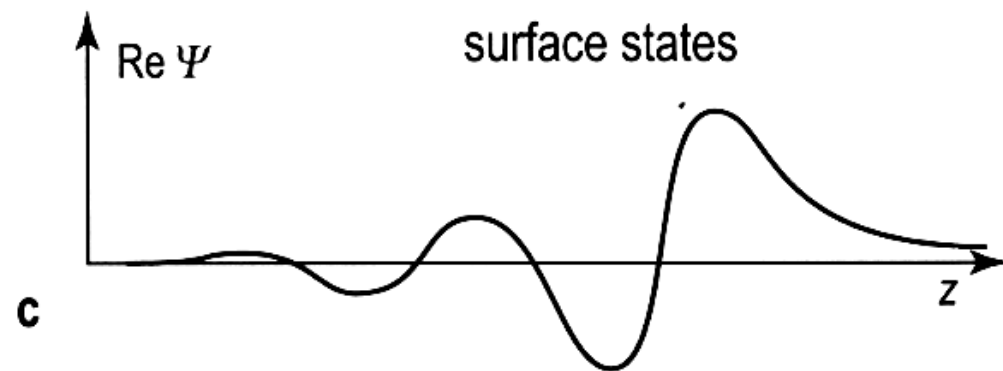
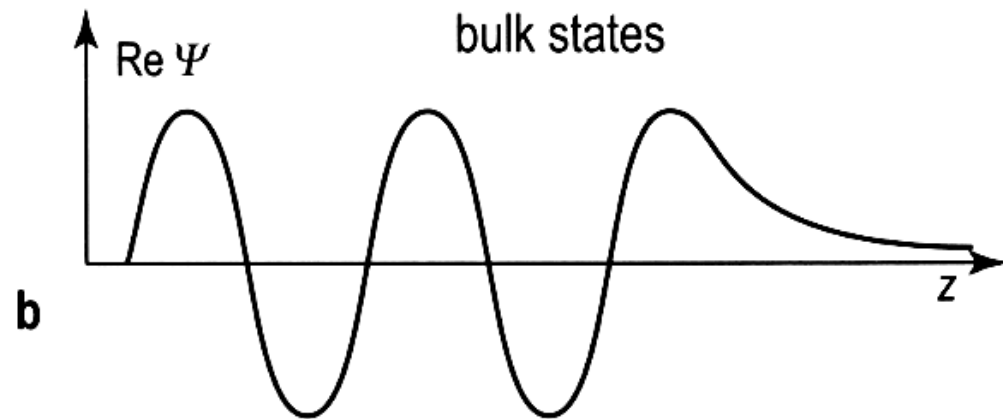
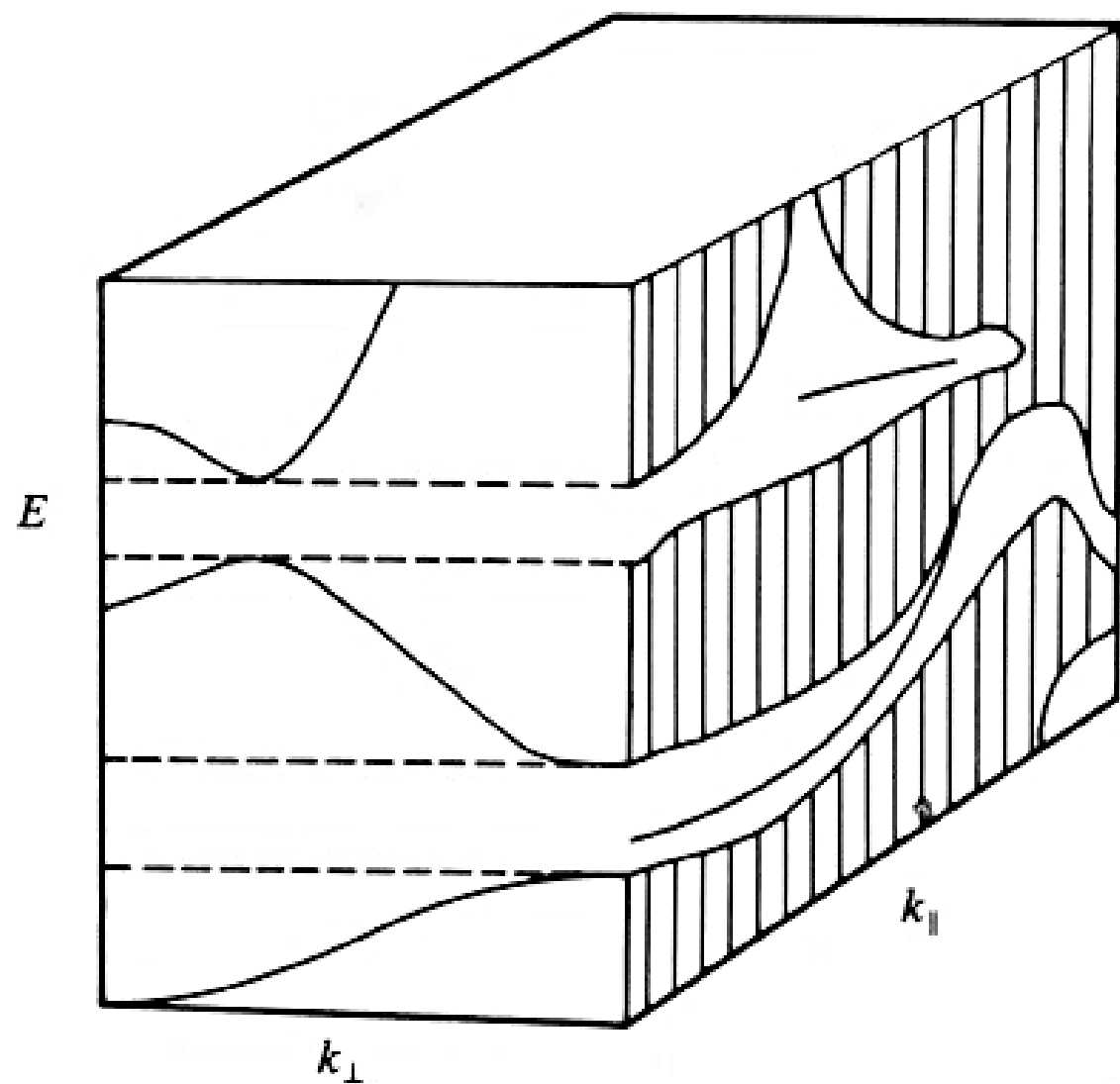
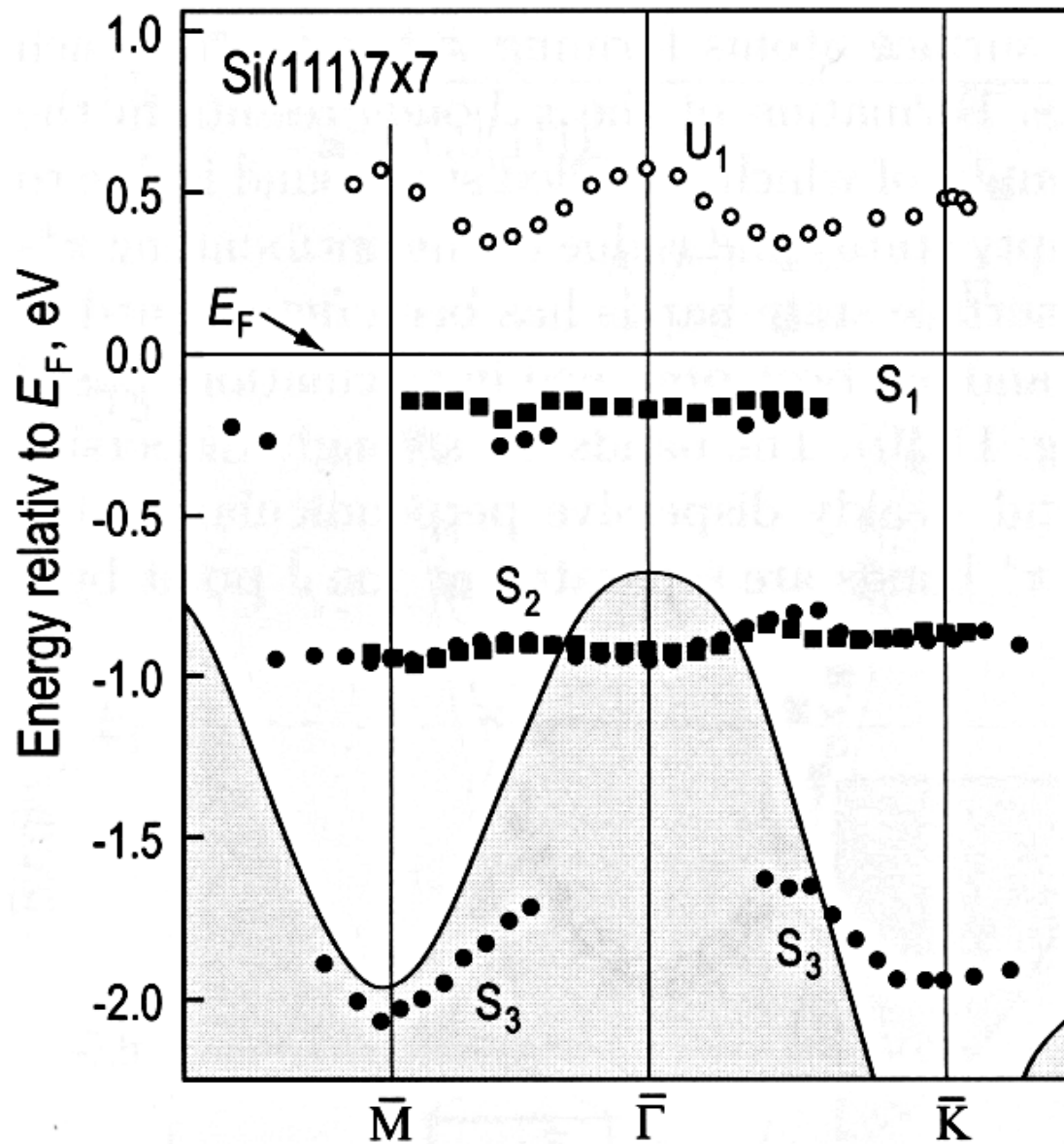
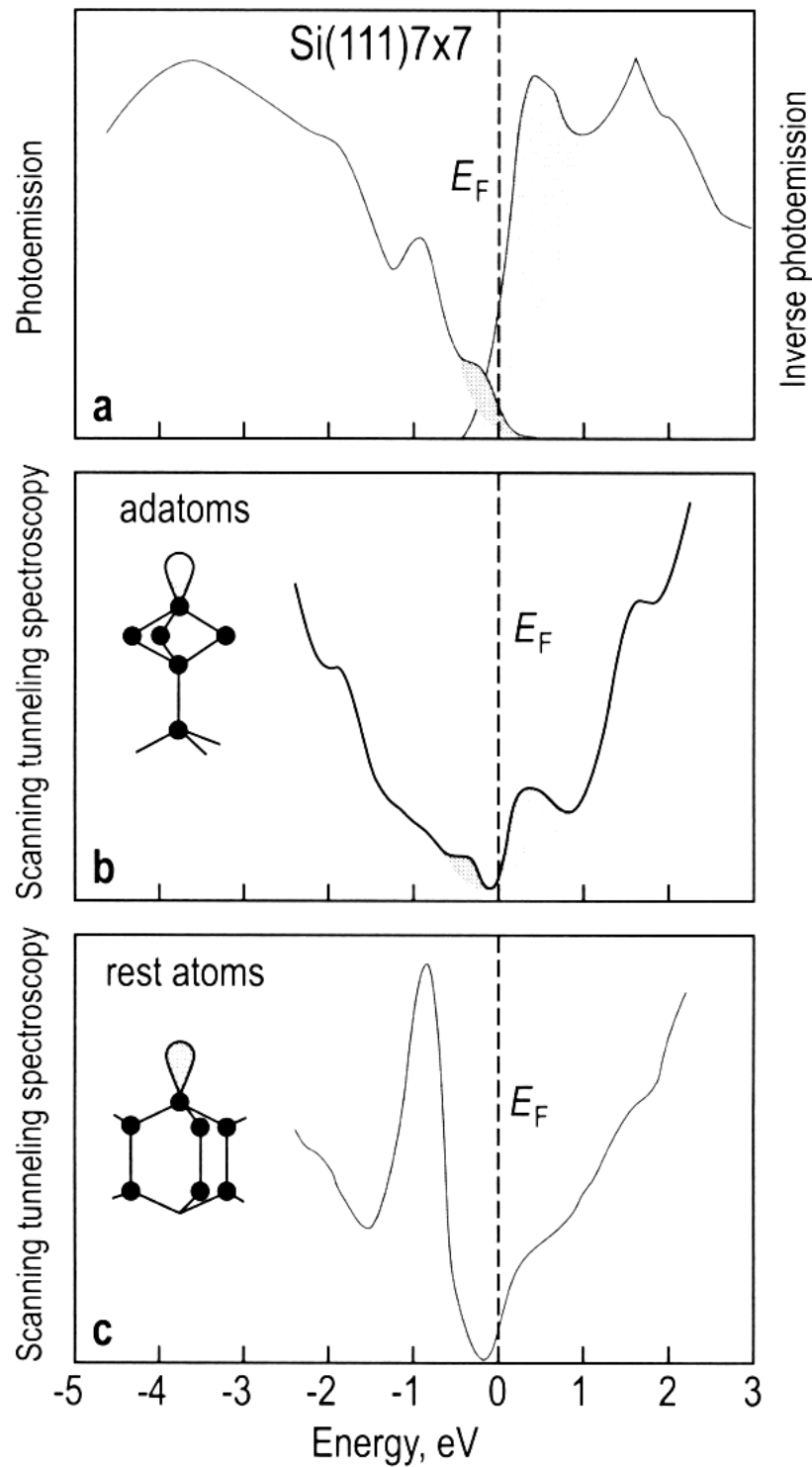


Fig. 4.15. Projected bulk band structure at the surface of a metal. The dispersion of two possible surface state bands is indicated.

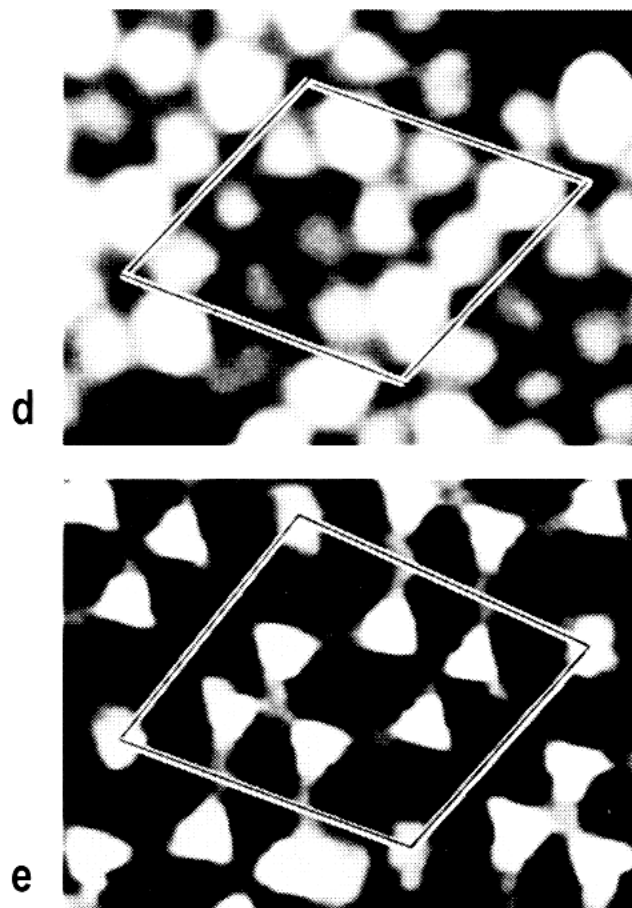




**Fig. 11.11.** Surface band structure of the Si(111)7 $\times$ 7 surface (referring to the Brillouin zone for the Si(111)1 $\times$ 1 surface) as resulting from ARUPS (closed circles [11.12] and closed squares [11.13]) and KRIPES (open circles [11.12]) measurements. The shaded region is the projected bulk band structure

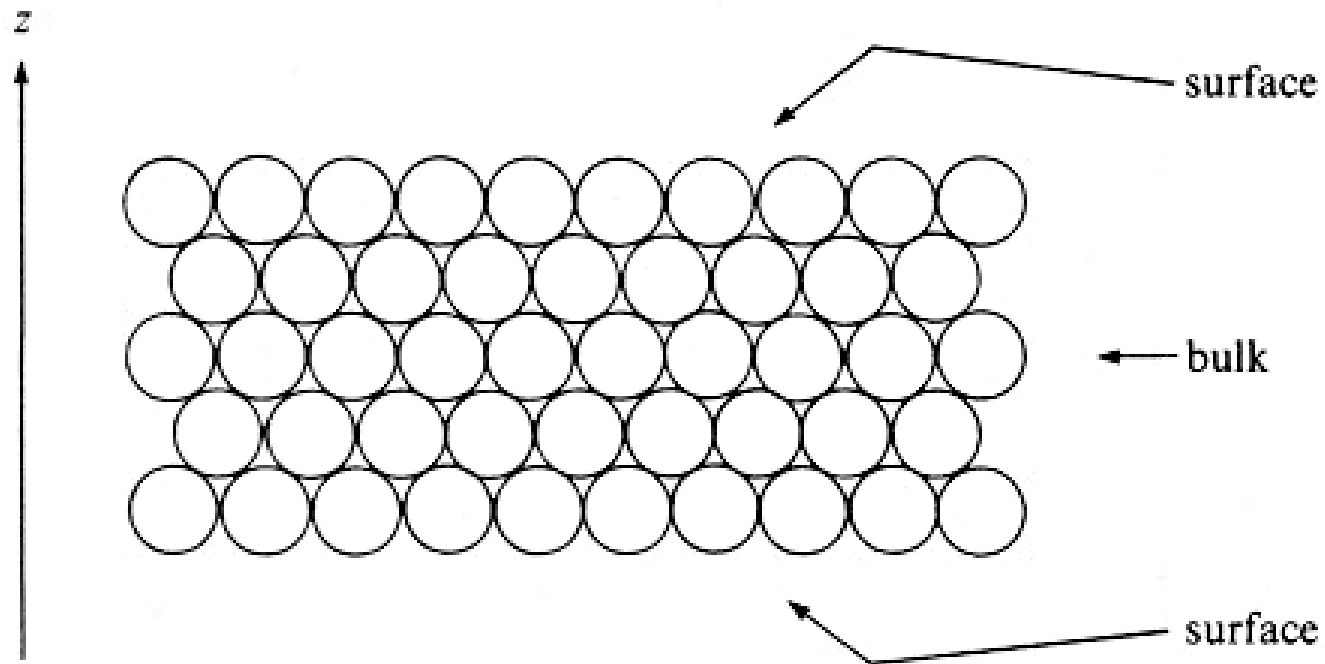


**Fig. 11.12.** Surface states on the Si(111)7×7 surface detected with (a) photoemission, inverse photoemission [11.15] and b, (c) scanning tunneling spectroscopy [11.16]. Two types of dangling bond states occur; a lone pair at the rest atom sites and a partially filled state at the adatom sites. The surface localization of the states is revealed in STM/STS experiments by using the bias voltage corresponding to the energy position of a given state, (d) at  $-0.35$  V for the adatom state and (e) at  $-0.8$  V for the rest atom state [11.14] (after Himpsel [11.17])



# Slab model

Fig. 4.16. Edge view of a five-layer slab for surface electronic structure calculations. Slab has infinite extent in the  $x$ - $y$  plane.



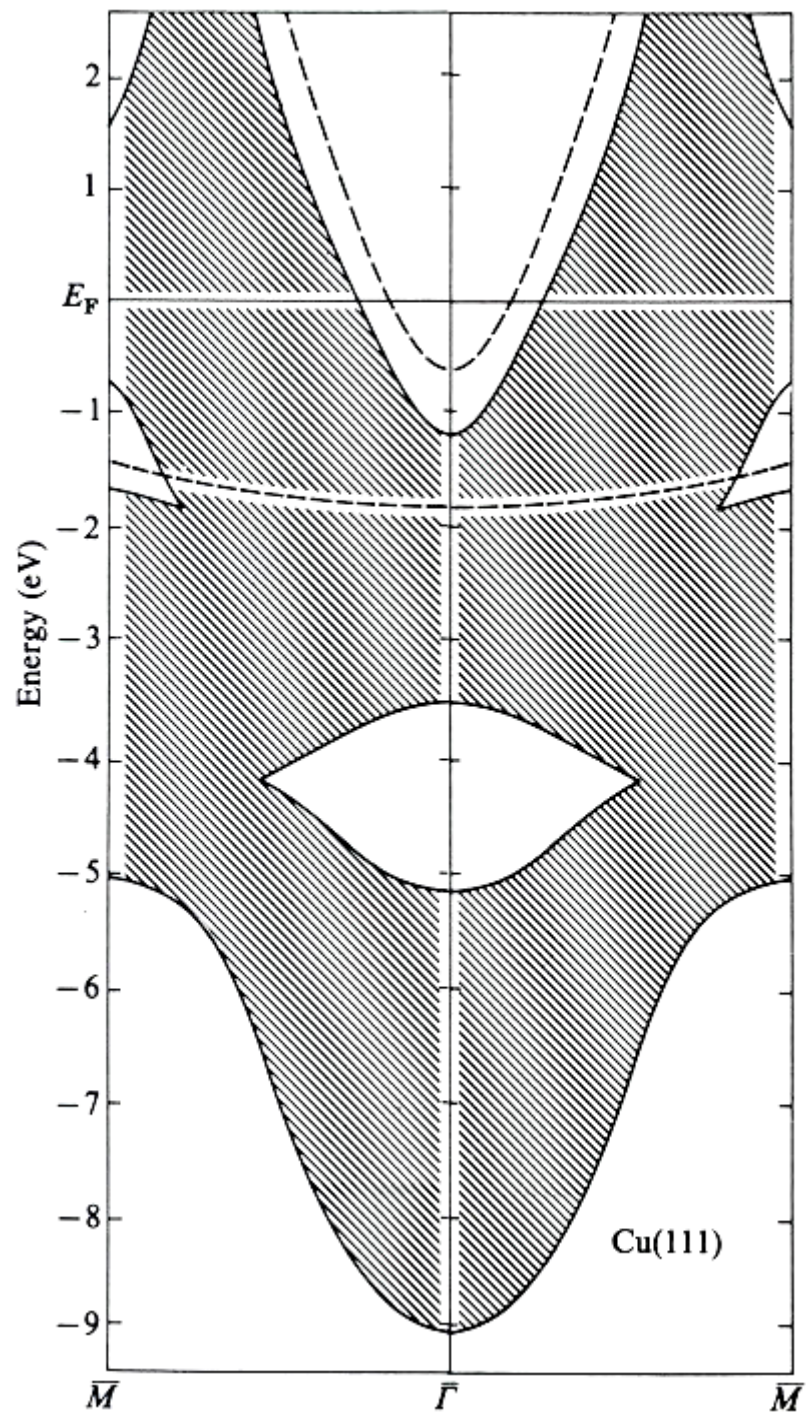
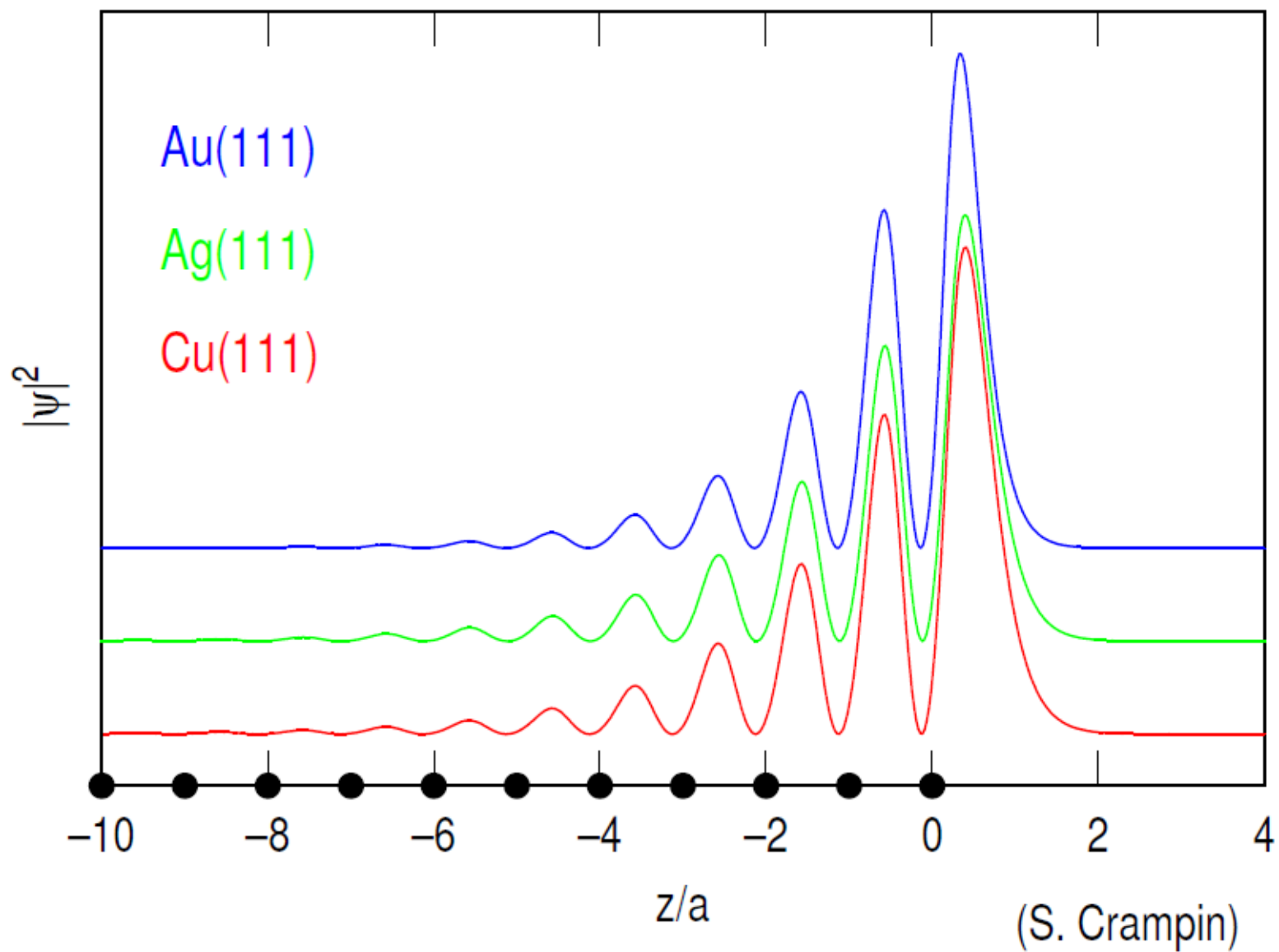


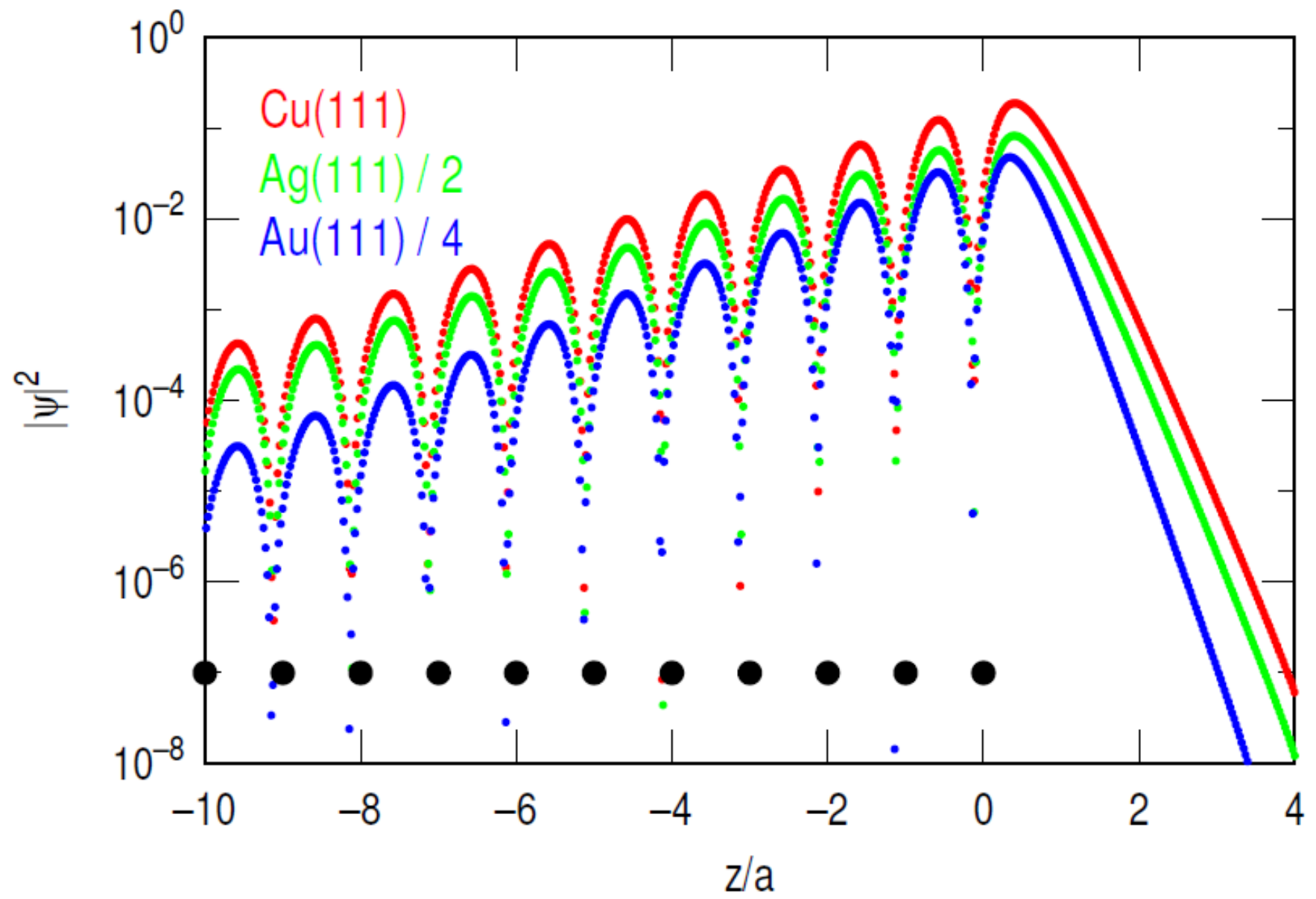
Fig. 4.17. Surface states (dashed curves) and bulk projected bands at a Cu(111) surface according to a six-layer surface band structure calculation (Euceda, Bylander & Kleinman, 1983).



# SURFACE STATE WAVEFUNCTIONS

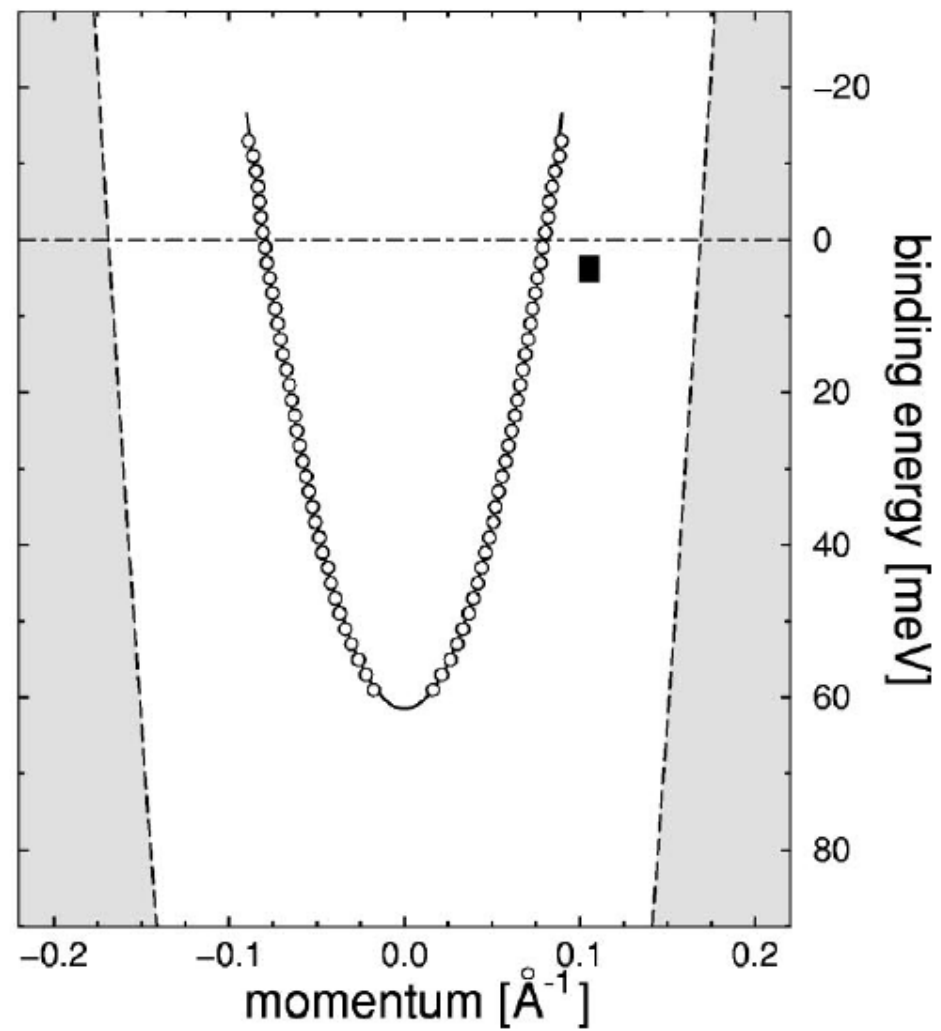
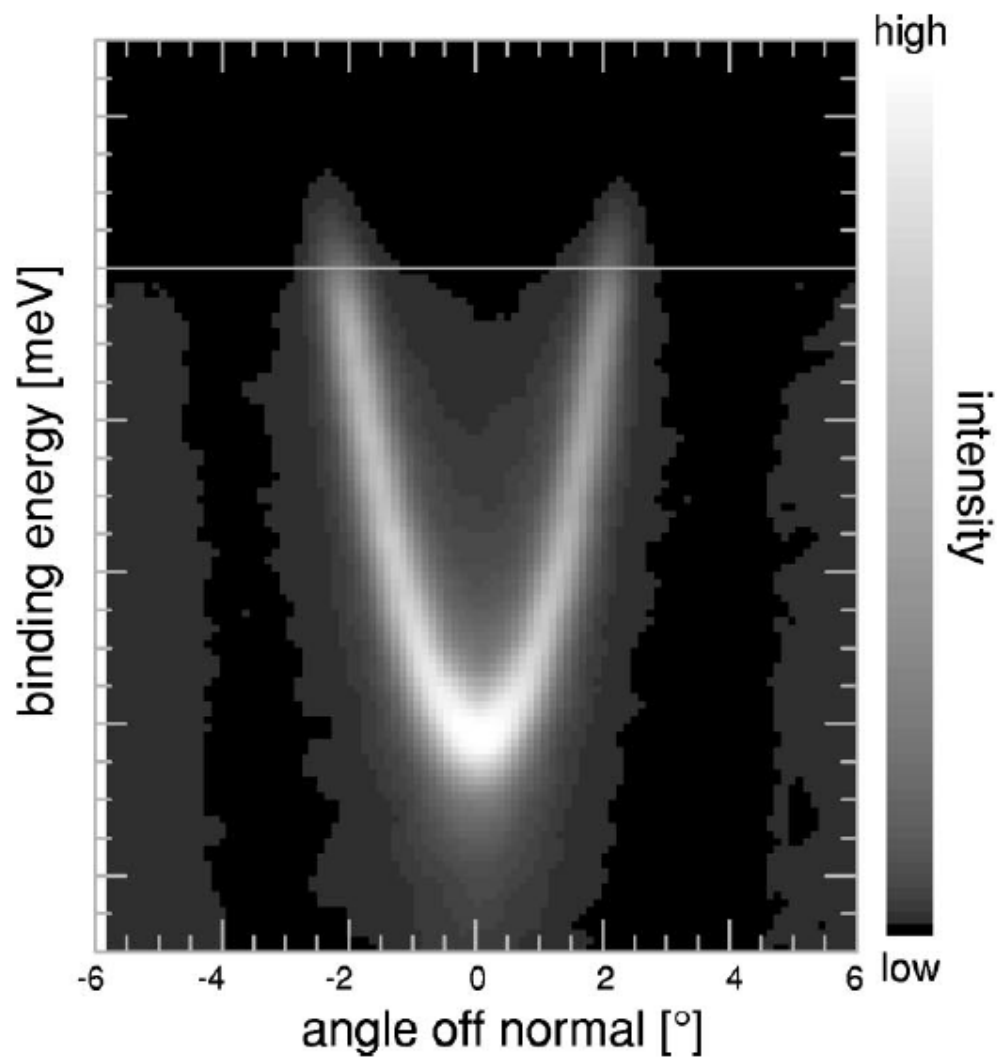


# SURFACE STATE WAVEFUNCTIONS

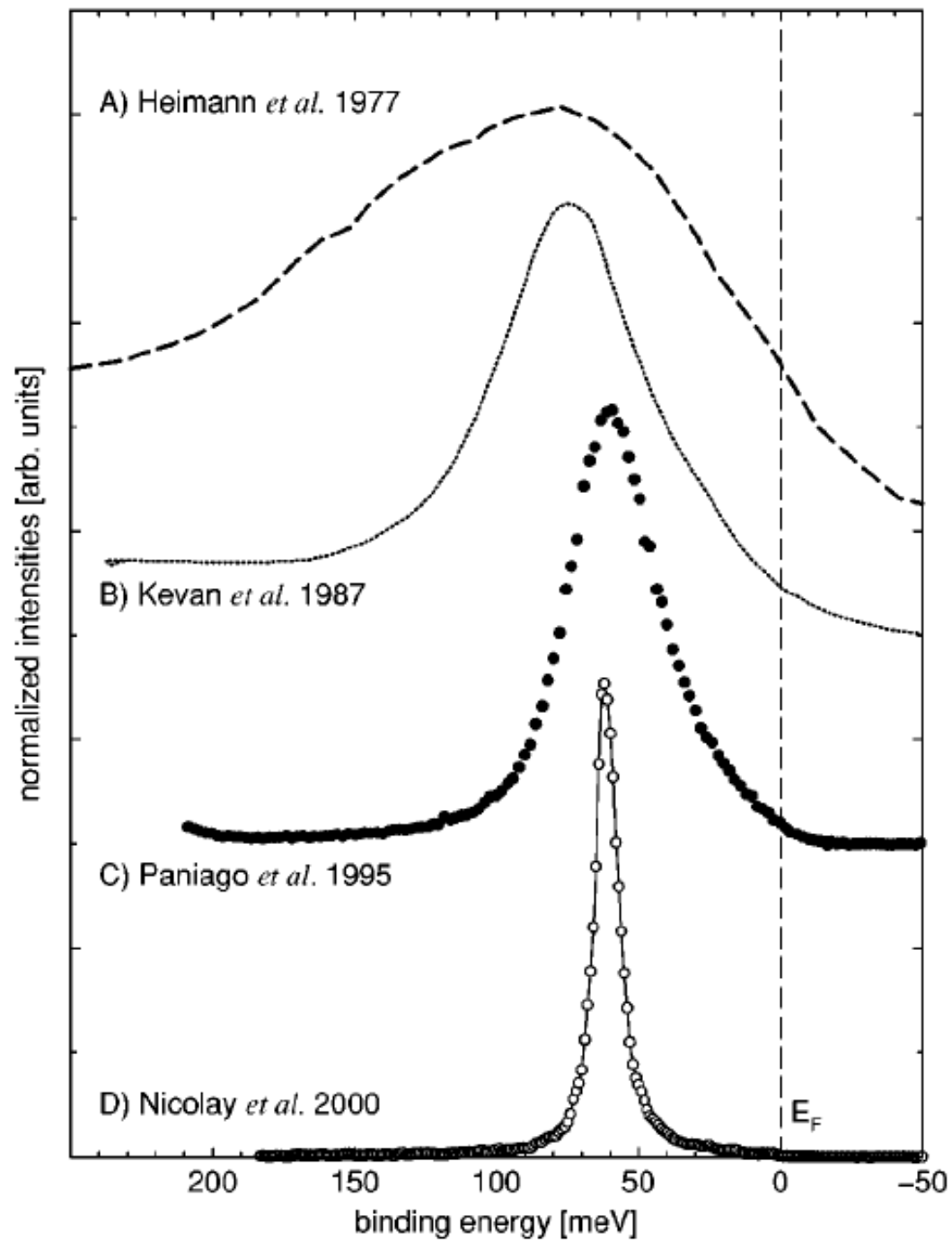


(S. Crampin)

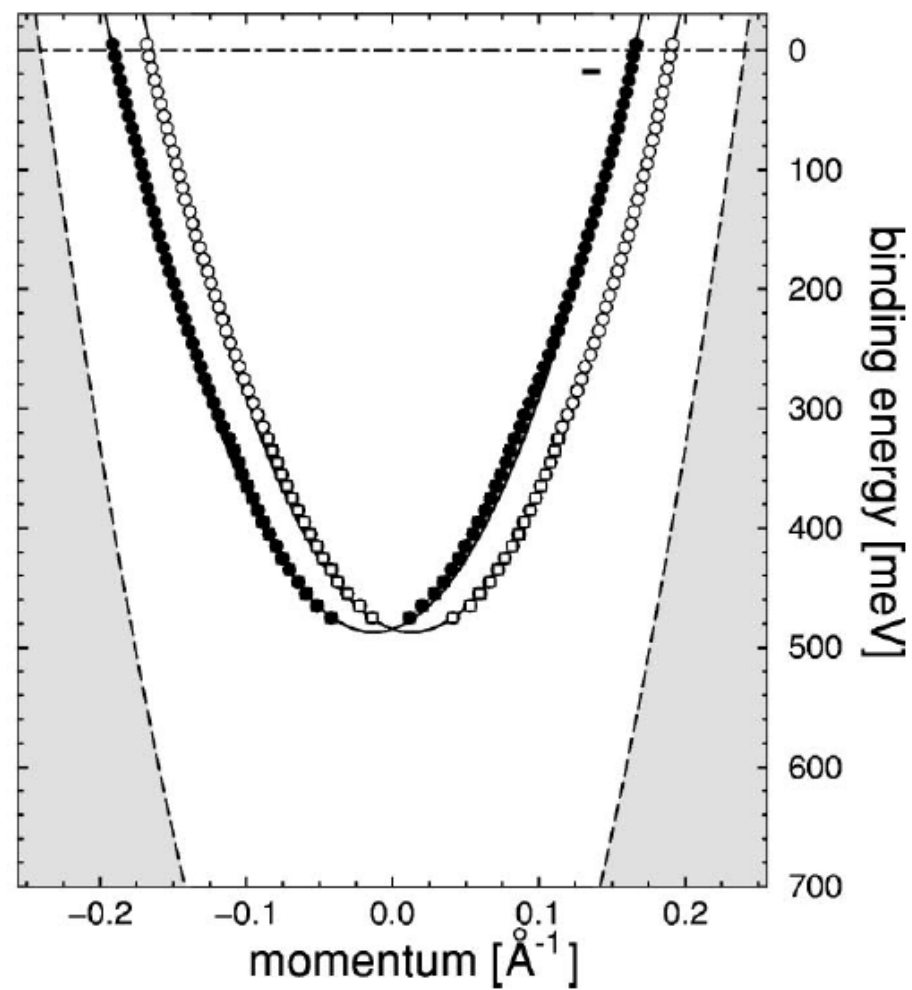
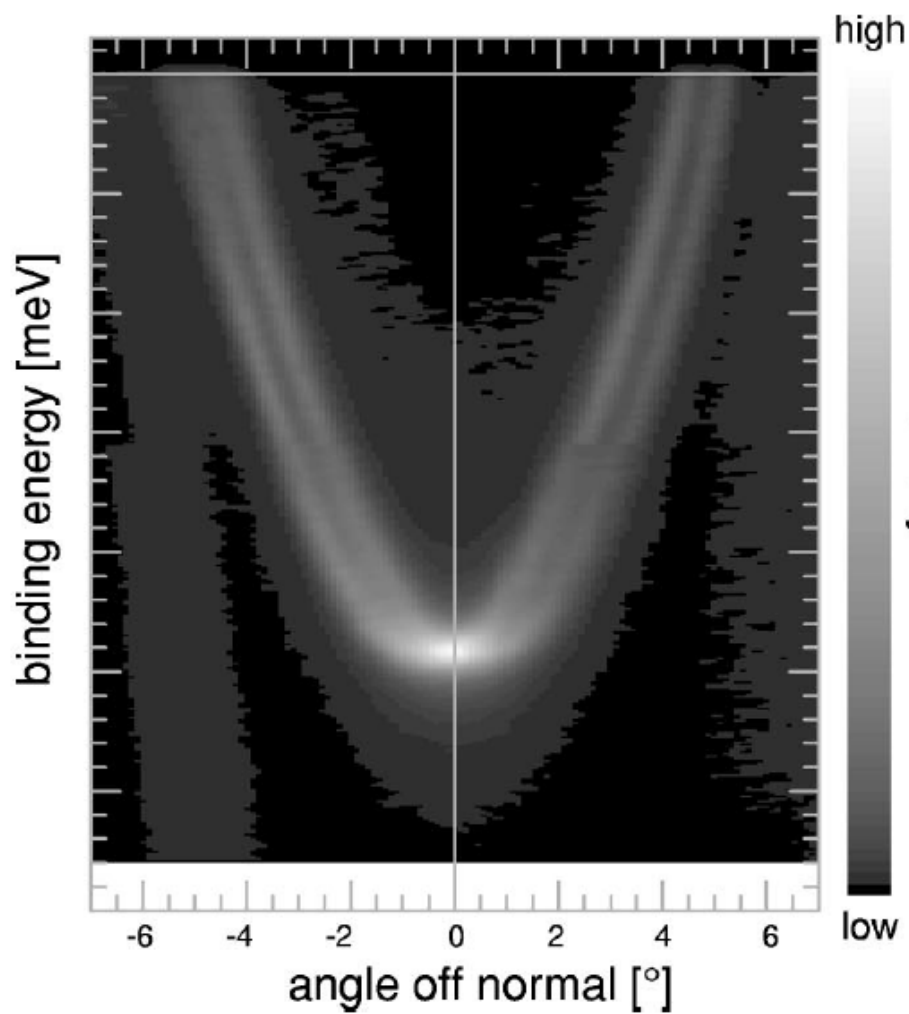
# Ag(111) surface state



# Ag(111) L-Gap Surface State by PES



# Au(111) surface state



relativistic correction to Schrödinger equation:

PRL **77**, 3419 (1996)

$$H_{\text{SOC}} = (\hbar/4mc^2) (\nabla V \times \vec{p}) \cdot \vec{\sigma}$$

$V$ : external potential,  $p$ : momentum,  $\sigma$ : Pauli spin operator

approximately proportional to:  $\vec{L} \cdot \vec{S}$  **SOC**

→ spin degenerate levels may be spin split by SOC

but: forbidden for systems inversion symmetry

crystal surface breaks this symmetry!

Au(111):

Fermi surfaces and (one half of) the surface Brillouin zone.

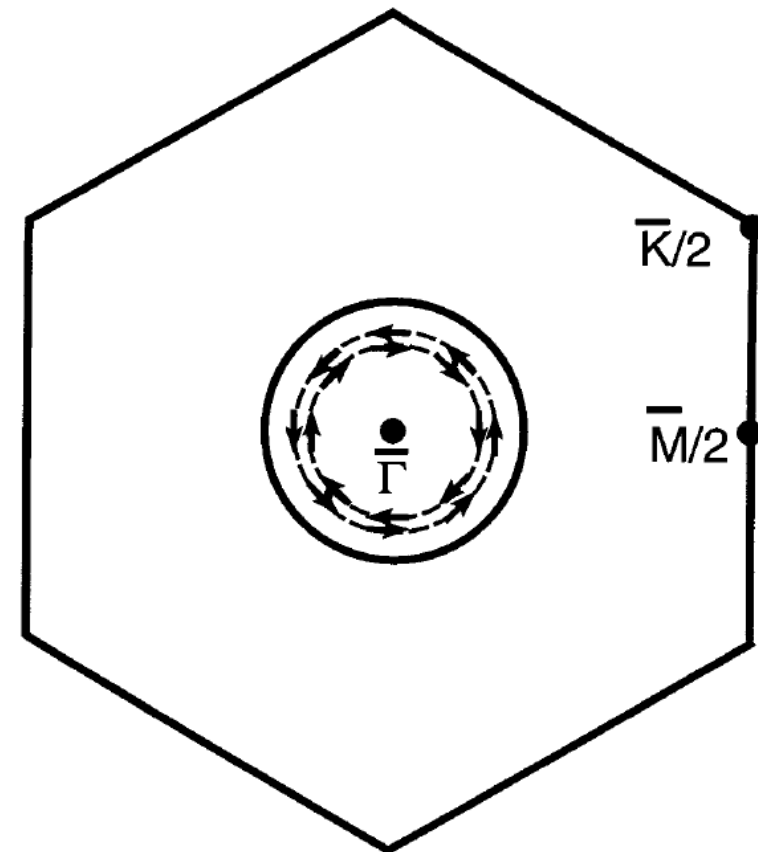
Arrows and dashed lines indicate spin orientations

Solid line represents the bulk Fermi surface neck.

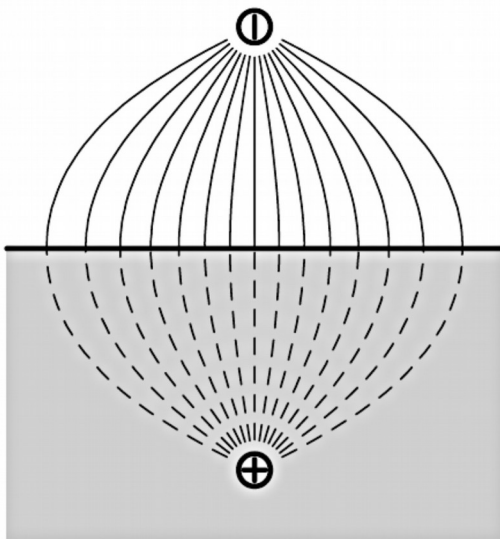
All spins are in the plane of the surface.

At the surface:  $k_F = -0.153, 0.177 \text{ \AA}^{-1}$ ,  $M = 1.26 \text{ \AA}^{-1}$ .

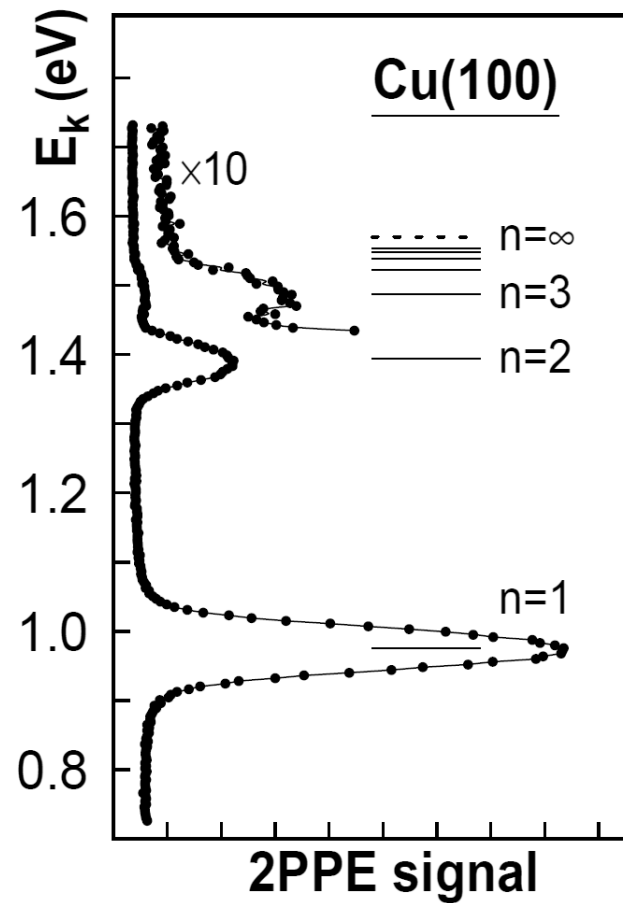
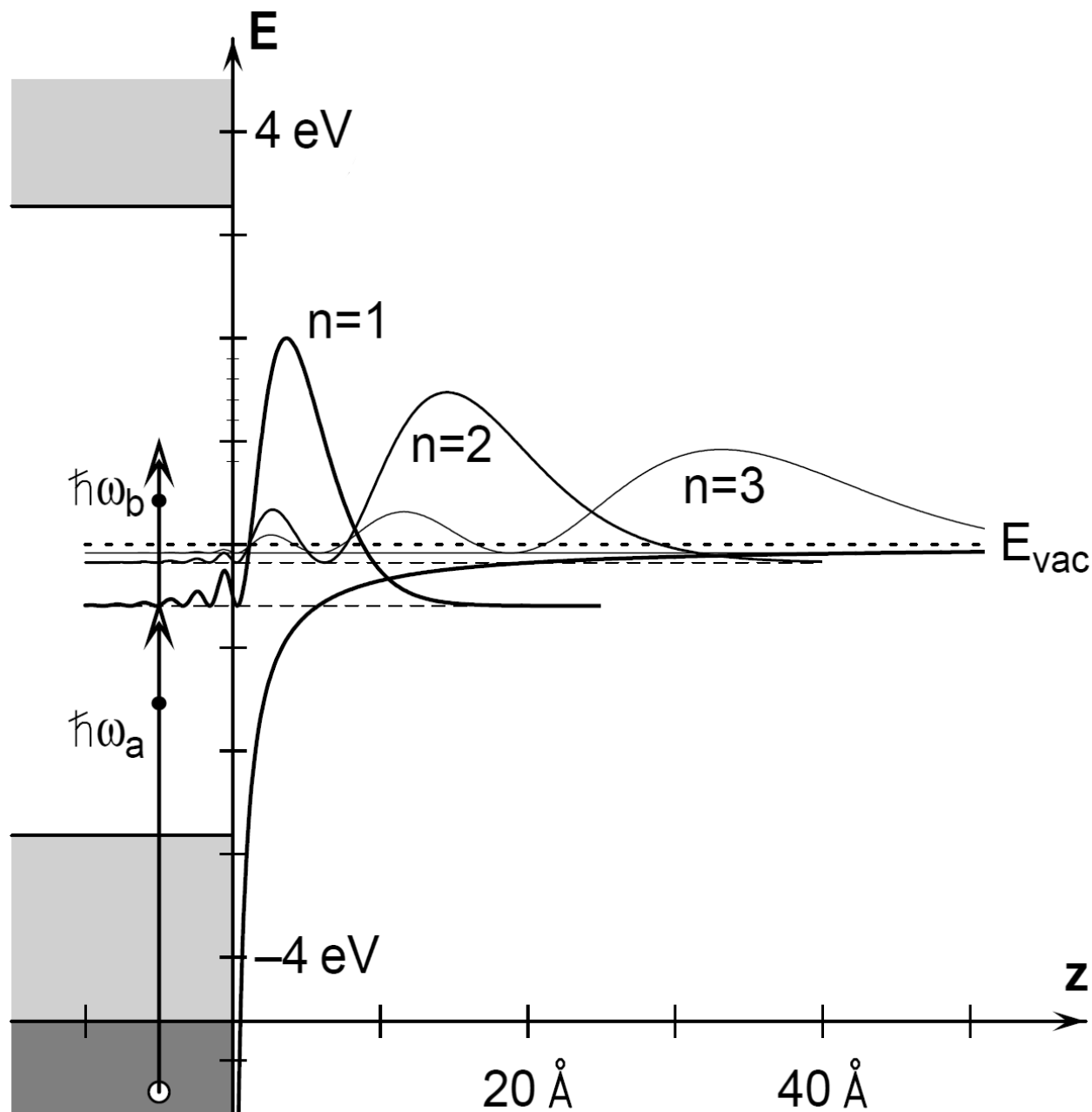
In the bulk:  $k_F = -0.216 \text{ \AA}^{-1}$ .



# Bildpotentialzustände



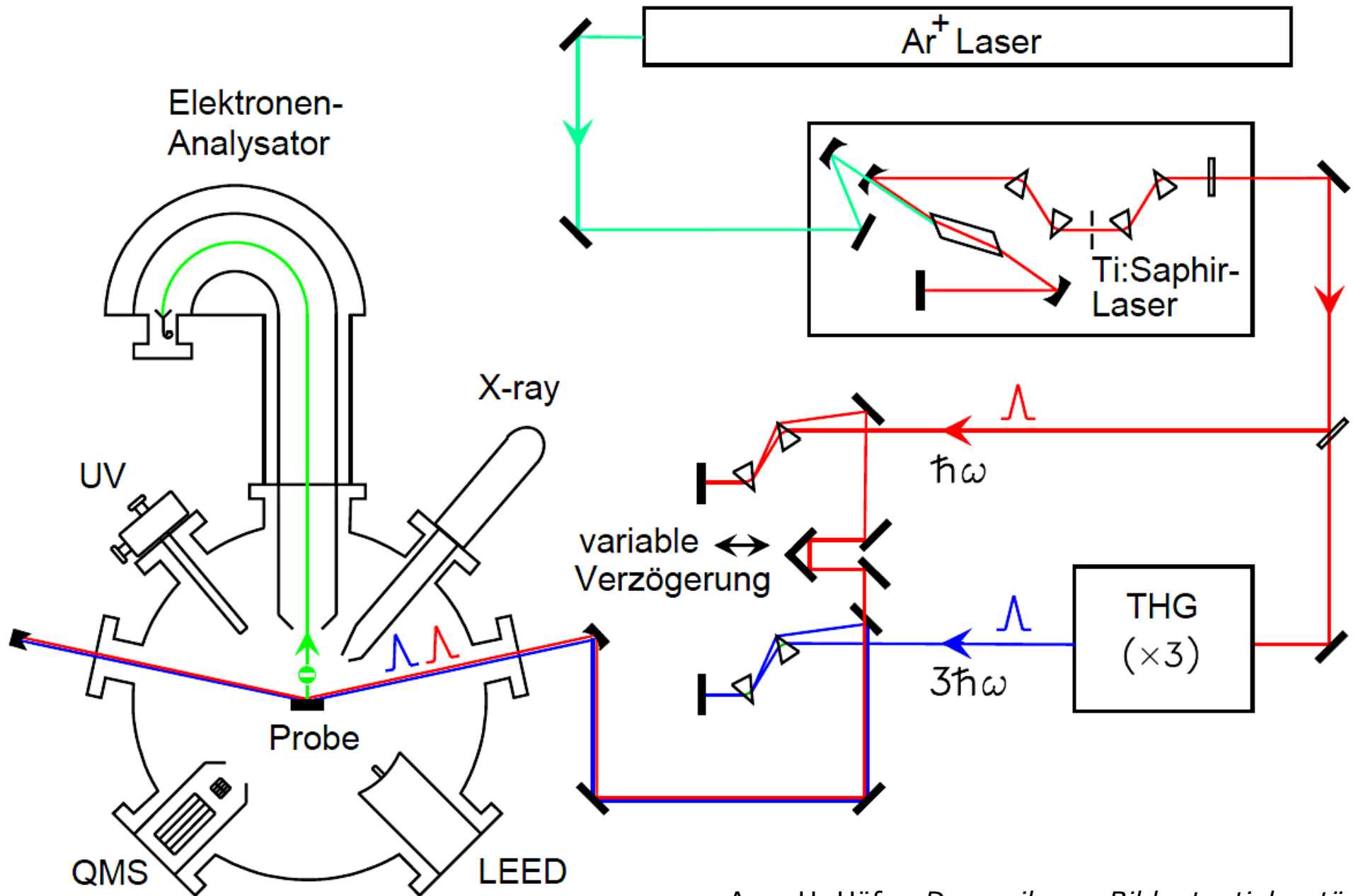
$$E_n = E_{\text{vac}} - \frac{0.85 \text{ eV}}{(n + a)^2}, \quad n = 1, 2, \dots$$

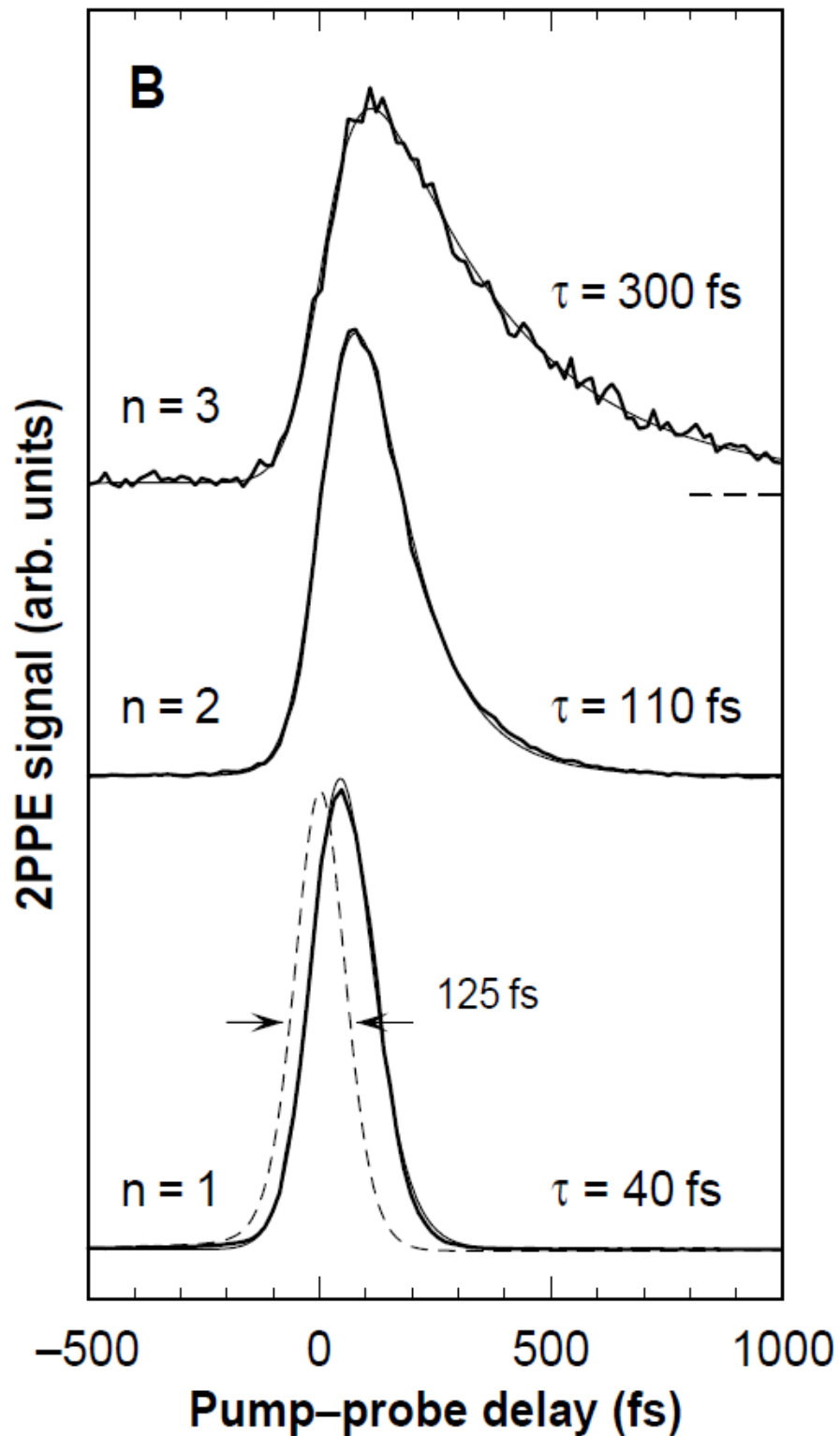


Photoelektronen-Spektrum  
 Aufgrund sequentieller  
 Anregung mit Photonen der  
 Energien  $h\nu_a$  und  $h\nu_b$



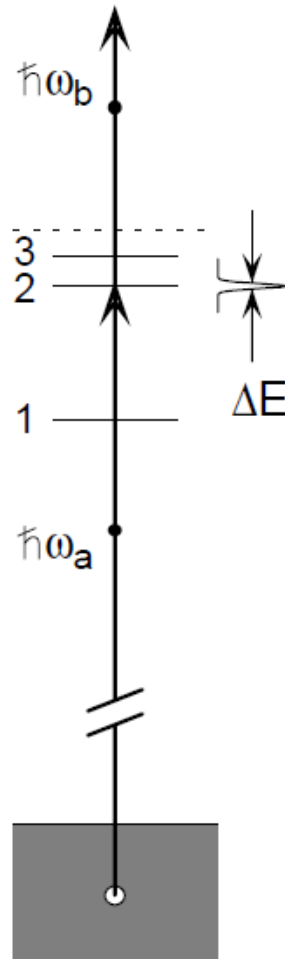
# Zweiphotonenphotoemission (2PPE)





Cu(100)

$E_B > 80$  meV



Zeitabhängigkeit des 2PPE-Signals  
der niedrigsten drei Bildpotentialzustände  
von Cu(100).

Gestrichelte Linie: Kreuzkorrelation der  
Pump- und Probeimpulse ohne  
resonante Zwischenzustände.

Die geringe spektrale Breite  $\Delta E$  der  
Pumpimpulse auf der Energieskala der  
Bildpotentialzustände verdeutlicht das  
Anregungsschema rechts.